

Incompressible and Inviscid Fluid Equations

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1 Introduction

In this file, we will develop the fluids equation for conservation of mass or *Continuity Equation* and special form of the conservation of momentum or *Euler Equation*.

2 The Continuity Equation

In this section we will derive the continuity equation. For this derivation we will be working with some volume of water, which we will call D . This could be a lake, the ocean, or some more controlled system such as a tank. Later when we begin our work with waves, we will specify exactly what we want D to be.

Let $\mathbf{u} = (u_1, u_2, u_3)$ be a velocity field, where $u_i = u_i(x, y, z, t)$ for $i = 1, 2, 3$.

The density of the water at a point (x, y, z) and at time t will be denoted as $\rho(x, y, z, t)$.

Let V be some arbitrary control bounded volume contained in D , and S the piecewise smooth and closed surface of V .

The mass of V can be written as $m = \int_V \rho(x, y, z, t) dV$.

The rate of change in mass for the volume V can then be written as $d/dt (\int_V \rho dV)$.

The principle of conservation of mass established that the rate of change of mass inside V equals the net mass that passes through the surface S per unit of time. Then,

$$\int_V \frac{\partial \rho}{\partial t} = - \int_S \rho \mathbf{u} \cdot \mathbf{n} dS, \quad (2.1)$$

where $\mathbf{n}$ is the outward unit normal. A quick check shows that the units work out correctly. In fact, $\frac{kg}{s} = \frac{(kg)(m^3)}{(m^3)(s)} = \frac{(kg)(m)(m^2)}{(m^3)(s)}$.

The integral on the left is obtained by applying Leibniz's rule. It means

$$\frac{d}{dt} \int_V \rho dV = \int_V \frac{\partial \rho}{\partial t} dV$$

The conditions are that

$\rho(\mathbf{x}, t)$ is an integrable function on V for each $t \in \mathbb{R}^+$ and

$\frac{\partial \rho}{\partial t}(\mathbf{x}, t)$ exists and is continuous in $V \times \mathbb{R}^+$.

Remark 2.1. The negative in the equation (2.1) can be explained as follows: Adding mass to a system results in a positive change of mass per unit of time. This flow of mass occurs through the surface S in the direction of the normal. Since S points outward, it must be negated to reflect a positive inflow of mass.

The Divergence Theorem applied on (2.1) leads to

$$\begin{aligned} \int_V \frac{\partial \rho}{\partial t} dV &= - \int_S (\rho \mathbf{u}) \cdot \mathbf{n} dS = - \int_V \nabla \cdot (\rho \mathbf{u}) dV \\ \int_V \left[\frac{\partial \rho}{\partial t} + \nabla \cdot (\rho \mathbf{u}) \right] dV &= 0. \end{aligned} \quad (2.2)$$

Since the volume V is arbitrary, then the integrand should be zero, *i. e.*,

$$\frac{\partial \rho}{\partial t} + \nabla \cdot (\rho \mathbf{u}) = 0, \quad (2.3)$$

which is known as the continuity equation.

Question 1. If the fluid is incompressible, then the density is constant. Show that the continuity equation reduces to

$$\nabla \cdot \mathbf{u} = 0. \quad (2.4)$$

Assuming constant density for water is a common assumption, which we will make for the remainder of this paper.

3 Euler's Equation of Motion

Now, we will derive Euler's equation of motion. Now, we will use the principle of *Conservation of Momentum*. We are still considering the same domain D , velocity field $\mathbf{u}$, and volume V with a surface S .

Previous to the derivation, we will proof some lemmas that will be used later.

Question 2. Prove the following three lemmas:

Lemma 3.1. If $f(x, y, z)$ is a scalar differentiable function in V and $\mathbf{u}(x, y, z)$ is a differentiable vector function in V , then

$$\nabla \cdot (f\mathbf{u}) = \nabla f \cdot \mathbf{u} + f \nabla \cdot (\mathbf{u})$$

Prove this lemma introducing the components of $f\mathbf{u}$.

Lemma 3.2. If $\mathbf{u} = \langle u_1, u_2, u_3 \rangle$ is a differentiable vector function, then

$$(\mathbf{u} \cdot \nabla) \mathbf{u} = \langle \mathbf{u} \cdot \nabla u_1, \mathbf{u} \cdot \nabla u_2, \mathbf{u} \cdot \nabla u_3 \rangle$$

Lemma 3.3. If f is a continuously differentiable scalar function, V a bounded region in the space, and S the surface enclosing the region V , then

$$\int_S f \mathbf{n} dS = \int_V \nabla f dV$$

where $\mathbf{n}$ is the outward unit normal.

To derive Euler's equation, we will consider the law of conservation of momentum.

The momentum, I inside the volume V , can be written as $I = \int_V (\rho \mathbf{u}) dV$,

then the rate of change of momentum is given by $dI/dt = d/dt \int_V (\rho \mathbf{u}) dV$.

Also the net inflow of momentum in V through the surface S is given by $-\int_S (\rho \mathbf{u}) \mathbf{u} \cdot \mathbf{n} dS$.

By the Law of Conservation of Momentum, the rate of change of momentum within the volume V equals the net inflow of momentum through the surface S plus the force $\mathbf{F}_p$ due to the pressure forces acting on the surface S and the force $\mathbf{F}_e$ due to the body forces per unit of mass acting inside the volume V . Thus we have

$$\frac{dI}{dt} = \frac{d}{dt} \int_V (\rho \mathbf{u}) dV = - \int_S (\rho \mathbf{u}) \mathbf{u} \cdot \mathbf{n} dS + \mathbf{F}, \quad (3.1)$$

where $\mathbf{F} = \mathbf{F}_p + \mathbf{F}_e$. A quick check insures that the units work out correctly, $\frac{(kg)(m)}{(s^2)} = \frac{(kg)(m)(m^3)}{(s)(m^3)(s)} = \frac{(kg)(m)(m)(m^2)}{(m^3)(s)(s)}$.

Upon expanding (3.1) component wise, we get

$$\int_V \frac{\partial}{\partial t} (\rho \mathbf{u}) dV = - \int_S \langle (\rho u_1 \mathbf{u}) \cdot \mathbf{n}, (\rho u_2 \mathbf{u}) \cdot \mathbf{n}, (\rho u_3 \mathbf{u}) \cdot \mathbf{n} \rangle dV + \mathbf{F}.$$

Question 3. Using Green's theorem, the incompressibility condition and the above lemmas show that

$$\int_V \rho \left[\frac{\partial \mathbf{u}}{\partial t} + (\mathbf{u} \cdot \nabla) \mathbf{u} \right] dV = \mathbf{F}. \quad (3.2)$$

Notice that $\mathbf{F}$ is the total force acting on the fluid, so we can break it up so that $\mathbf{F} = \mathbf{F}_p + \mathbf{F}_e$, where $\mathbf{F}_p$ is the force due to pressure, and $\mathbf{F}_e$ is some other external force. Pressure is force per unit area, so we can write

$$\mathbf{F}_p = - \int_S P \mathbf{n} dS \quad (3.3)$$

where $P = P(x, y, z, t)$ is the pressure. Again note the negative sign since the unit normal is outward. Also, if $\mathbf{A}_e$ represents the body forces per unit of mass, then the external force $\mathbf{F}_e$ is given by

$$\mathbf{F}_e = \int_V \rho \mathbf{A}_e dV \quad (3.4)$$

By substituting (3.3), and (3.4) into (3.2) and applying the Lemma 3.3 arrive to the final Euler's equation of motion given by

$$\frac{\partial \mathbf{u}}{\partial t} + (\mathbf{u} \cdot \nabla) \mathbf{u} = \mathbf{A}_e - \frac{\nabla P}{\rho}. \quad (3.5)$$